

FOURIER TRANSFORM AND HEAT FLOW

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We can extend the Fourier series solution for the heat flow problem to an infinite rod by using the Fourier transform. The solution satisfies the heat equation

$$\frac{\partial T(x,t)}{\partial t} = \frac{\partial^2 T(x,t)}{\partial x^2} \quad (1)$$

with the initial condition

$$T(x,0) = f(x) \quad (2)$$

The Fourier transform of $f(x)$ is

$$G(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx \quad (3)$$

with the inverse transform

$$f(x) = \int_{-\infty}^{\infty} G(\omega) e^{i\omega x} d\omega \quad (4)$$

A solution to 1 is given by

$$T(x,t) = \int_{-\infty}^{\infty} G(\omega) e^{i\omega x} e^{-\omega^2 t} d\omega \quad (5)$$

To verify this, we have, assuming we can differentiate under the integral sign

$$\frac{\partial T(x,t)}{\partial t} = \int_{-\infty}^{\infty} G(\omega) e^{i\omega x} (-\omega^2) e^{-\omega^2 t} d\omega \quad (6)$$

$$\frac{\partial^2 T(x,t)}{\partial x^2} = \int_{-\infty}^{\infty} G(\omega) (i\omega)^2 e^{i\omega x} e^{-\omega^2 t} d\omega \quad (7)$$

$$= \int_{-\infty}^{\infty} G(\omega) e^{i\omega x} (-\omega^2) e^{-\omega^2 t} d\omega \quad (8)$$

$$= \frac{\partial T(x,t)}{\partial t} \quad (9)$$

The initial condition 2 is

$$T(x,0) = \int_{-\infty}^{\infty} G(\omega) e^{i\omega x} d\omega = f(x) \quad (10)$$

We can eliminate $G(\omega)$ from 5 by using 3.

$$T(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega \int_{-\infty}^{\infty} d\xi f(\xi) e^{-i\xi\omega} e^{i\omega x} e^{-\omega^2 t} \quad (11)$$

where we've introduced the dummy variable ξ for the second integration. We can reverse the order of integration to get

$$T(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\xi f(\xi) \int_{-\infty}^{\infty} d\omega e^{-\omega^2 t + i\omega(x-\xi)} \quad (12)$$

The ω integral can be evaluated explicitly by completing the square in the exponent. The exponent is

$$-\omega^2 t + i\omega(x-\xi) = -t \left(\omega^2 - \frac{i\omega(x-\xi)}{t} \right) \quad (13)$$

$$= -t \left(\omega - \frac{i(x-\xi)}{2t} \right)^2 - t \frac{(x-\xi)^2}{4t^2} \quad (14)$$

$$= -t \left(\omega - \frac{i(x-\xi)}{2t} \right)^2 - \frac{(x-\xi)^2}{4t} \quad (15)$$

The last term doesn't involve ω so can be pulled outside the ω integral, giving

$$T(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\xi f(\xi) e^{-(x-\xi)^2/4t} \int_{-\infty}^{\infty} e^{-t(\omega - i(x-\xi)/2t)^2} d\omega \quad (16)$$

The last integral can be evaluated using the substitution

$$u = \sqrt{t}(\omega - i(x-\xi)/2t) \quad (17)$$

$$du = \sqrt{t} d\omega \quad (18)$$

$$\omega = -\infty \rightarrow u = -\infty \quad (19)$$

$$\omega = \infty \rightarrow u = \infty \quad (20)$$

The integral is then

$$\int_{-\infty}^{\infty} e^{-t(\omega - i(x-\xi)/2t)^2} d\omega = \frac{1}{\sqrt{t}} \int_{-\infty}^{\infty} e^{-u^2} du \quad (21)$$

$$= \sqrt{\frac{\pi}{t}} \quad (22)$$

where we used the result for the Gaussian integral

$$\int_{-\infty}^{\infty} e^{-u^2} du = \sqrt{\pi} \quad (23)$$

Inserting back into 16 we get

$$T(x, t) = \frac{1}{2\pi} \sqrt{\frac{\pi}{t}} \int_{-\infty}^{\infty} d\xi f(\xi) e^{-(x-\xi)^2/4t} \quad (24)$$

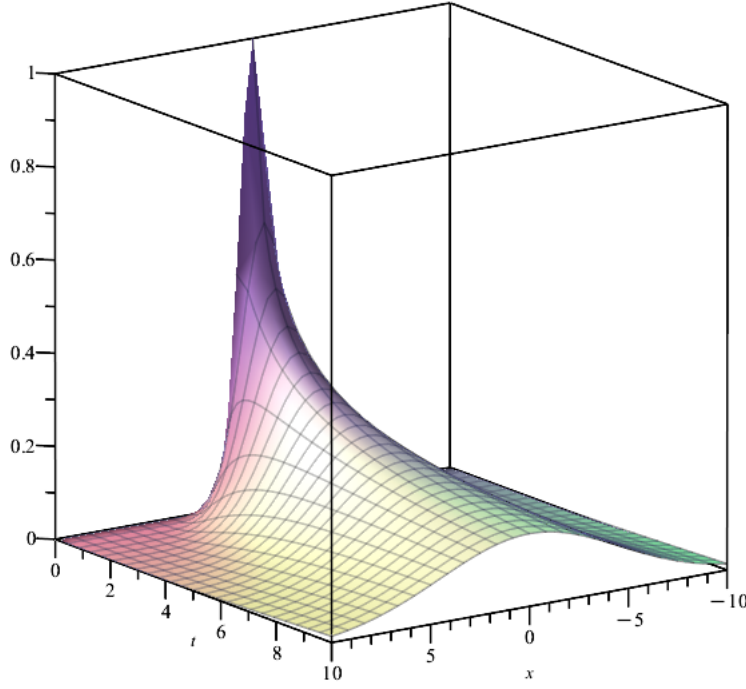


FIGURE 1. Plot of $T(x, t) = \frac{e^{-x^2/(4t+1)}}{\sqrt{4t+1}}$ for $x \in [-10, 10]$ and $t \in [0, 10]$.

$$= \frac{1}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} d\xi f(\xi) e^{-(x-\xi)^2/4t} \quad (25)$$

As an example, we take the initial condition to be

$$f(\xi) = e^{-\xi^2} \quad (26)$$

This gives

$$T(x, t) = \frac{1}{2\sqrt{\pi t}} \int_{-\infty}^{\infty} d\xi e^{-\xi^2} e^{-(x-\xi)^2/4t} \quad (27)$$

$$= \frac{e^{-x^2/(4t+1)}}{\sqrt{4t+1}} \quad (28)$$

where I used Maple to do the integral. To do it by hand, you could employ the completing the square technique used above.

The plot of $T(x, t)$ is given in Fig. 1. As we might expect, the temperature tends to decrease and spread out along the rod.